# ON CONDITIONS FOR THE EXISTENCE OF PERIODIC SOLUTIONS OF SYSTEMS OF DIFFERENTIAL EQUATIONS WITH DISCONTINUOUS RIGHT-HAND Sides containing a small parameter 

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1. Certain preliminary remarks. Consider the system of differential equations of the form

$$
\begin{equation*}
\frac{d x}{d t}=X(x, t)+\mu f(x, t, \mu) \tag{1.1}
\end{equation*}
$$

where $x, X, f$ are $n$-dimensional vectors, and $\mu$ is a small parameter.
It is supposed that:
a) The functions $X$ and $f$ are defined and single-valued for all real values of $t$, for all values of $\mu$ lying in an interval $0 \leqslant \mu \leqslant \mu_{0}$, and for all $x$ lying in an $n$-dimensional domain $G$;
b) for all $x$ and $\mu$ in question, the functions $X$ and $f$ are continuous and possess the period $T$ :

$$
X(x, t+T) \equiv X(x, t), \quad f(x, t+T, \mu) \equiv f(x, t, \mu)
$$

c) the domain $G$ is divided, by means of continuous smooth surfaces, into a finite number of subdomains $G_{k}$, in each of which, including its boundary, the function $X$ possesses continuous second-order partial derivatives with respect to $x$, and the function $f$ possesses continuous first-order partial derivatives with respect to $x$ and $\mu$, for $0<\mu \leqslant \mu_{0}$;
d) on the surfaces separating the domains $G_{k}$ (surfaces which will henceforth be referred to as discontinuity surfaces) there may occur discontinuities of the first kind of the functions $X$ and $f$, or of their
first-order partial derivatives with respect to $x$ and $\mu$ as the case may be, or of the second-order partial derivatives of $X$ with respect to $x$;
e) the equation of the surface of discontinuity between the domains $G_{k}$ and $G_{k+1}$ is taken to be of the form

$$
\begin{equation*}
\varphi_{k}(x)=0 \tag{1.2}
\end{equation*}
$$

It is assumed that the functions $\phi_{k}$ possess continuous second-order partial derivatives on the portions of the surfaces (1.2) which actually lie in the domain $G$.

It is supposed also that the generating system

$$
\begin{equation*}
\frac{d x_{0}}{d t}=X\left(x_{0}, t\right) \tag{1.3}
\end{equation*}
$$

possesses a family of periodic solutions, depending on $l$ independent parameters

$$
\begin{equation*}
x_{0}=x_{0}\left(t, h_{1}, \ldots, h_{l}\right) \tag{1.4}
\end{equation*}
$$

that all integral curves of this family pass through each of the $G_{k}$, that at all points of the intersection of one of these curves and surface (1.2) the following condition holds:

$$
\begin{equation*}
\frac{\partial \varphi_{k}}{\partial x} \frac{d x_{0}}{d t} \neq 0 \quad\left(\frac{\partial \varphi_{k}}{\partial x} \equiv \operatorname{grad} \varphi_{k}\right) \tag{1.5}
\end{equation*}
$$

and that the domain $G$ contains an $n$-dimensional neighborhood of each point of the $(l+1)$ dimensional manifold (1.4). From the results obtained in [1,2] it follows that, under the hypotheses enumerated explicitly above, there exist periodic solutions of the system (1.1) in the neighborhood of certain solutions (1.4), and tending continuously to them as $\mu \rightarrow 0$. In order for this to happen, the values of the parameters $h_{1}, \ldots, h_{l}$ in the corresponding generating solutions must satisfy certain $l$ conditions:

$$
\begin{equation*}
P_{i}\left(h_{1}, \ldots, h_{l}\right)-0 \quad(i-1, \ldots, l) \tag{1.6}
\end{equation*}
$$

These conditions were obtained in [1] in the form of a determinant of order $n-l+1$, for whose actual construction there is required the knowledge of the point transformation effected by the solutions of the system (1.1). Meanwhile, in the case of equations with "smooth" righthand sides, an integral form of conditions (1.6) has been discovered [3].

In the present paper it is shown that this (integral) form of the necessary condition for the existence of periodic solutions may be extended to the case of equations with discontinuous right-hand sides.
2. Equations for the initial conditions. The derivation of the conditions for the existence of the periodic solutions of Equation (1.1) which are near to a solution of (1.4) may be carried out without employing the method of point transformations.

Suppose that the integral curves (1.4) pass in succession through the domains $G_{1}, G_{2}, \ldots, G_{m}$. Since the domain $G_{m}$ is closed, these curves must return again to the domain $G_{1}$. The equation of the boundary surface between $G_{m}$ and $G_{1}$ is

$$
\varphi_{m}(x)=0
$$

On each of the domains $G_{k}$ the conditions for the existence and uniqueness of the solutions of the system (1.3) are fulfilled. Consequently, on each of these domains one has a general solution of this equation, depending on an initial vector $C_{k}$ and on an initial instant of time $t_{0 k}$ :

$$
\begin{equation*}
x_{0 k}=x_{0 k}\left(t, t_{0 k}, C_{k i}\right) \quad(k=1, \ldots, m) \tag{2.1}
\end{equation*}
$$

where

$$
x_{0 k}\left(t_{0 k}, t_{0 k}, C_{k}\right)=C_{k}
$$

The integral curves (1.4) must, in $G_{k}$, coincide with some of the curves (2.1). On the surfaces of discontinuity the conditions of continuity and periodicity must be satisfied. If $\tau_{k}$ denotes the instants at which the integral curves intersect the surfaces of discontinuity, then we must have

$$
\begin{gather*}
x_{0 k}\left(\tau_{k}, t_{0 k}, C_{k}\right)-x_{0, k+1}\left(\tau_{k}, t_{0, k+1} C_{k+1}\right)=0 \quad(k=1, \ldots, m-1)  \tag{2.2}\\
c_{0 m}\left(\tau_{m}, t_{0 m}, C_{m}\right)-x_{01}\left(\tau_{m}-T, t_{01}, C_{1}\right)=0  \tag{2.3}\\
\varphi_{k}\left[x_{0 k}\left(\tau_{k}, t_{0 k}, C_{k}\right)\right]=0 \quad(k-1, \ldots, m) \tag{2.4}
\end{gather*}
$$

Equation (1.3) has a family of solutions (1.4), therefore the system (2.2) to (2.4) must also have an $l$ parameter family of solutions of the form

$$
\begin{equation*}
C_{k}=C_{k}\left(h_{1}, \ldots, h_{l}\right), \quad \tau_{k}=\tau_{k}\left(h_{1}, \ldots, h_{l}\right) \tag{2.5}
\end{equation*}
$$

Here, $t_{0 k}$ may be chosen arbitrarily in the interval $\tau_{k-1} \leqslant t_{0 k} \leqslant \tau_{k}$. on the other hand, it is easily seen that the general solution of Equation (1.1) in $G_{k}$ is given by

$$
\begin{equation*}
x_{k}=x_{k}\left(t, t_{0 k}, D_{k}\right)=x_{0 k}\left(t, t_{0 k}, D_{k}\right)+\mu y_{k}\left(t, t_{0 k}, D_{k}, \mu\right) \quad(k=1, \ldots, m) \tag{2.6}
\end{equation*}
$$

where $t_{0 k}$ has the same value as in (2.2) to (2.4).

The initial vector $D$, at the instant $r_{k}^{\prime}$ at which the integral curves return to the surfaces of discontinuity, may be determined from the equations

$$
\begin{align*}
& x_{0 k}\left(\tau_{k}{ }^{\prime}, t_{0 k}, D_{k}\right)+\mu y_{k}\left(\tau_{k}^{\prime}, t_{0 k}, D_{k}, \mu\right)-x_{0, k+1}\left(\tau_{k}^{\prime}, t_{0, k+1}, D_{k+1}\right)- \\
& \quad-\mu y_{k+1}\left(\tau_{k}{ }^{\prime}, t_{0, k+1}, D_{k \cdot 1}, \mu\right)=0 \quad(k=1, \ldots, m-1)  \tag{2.7}\\
& x_{0 m}\left(\tau_{m}{ }^{\prime}, t_{0 m}, D_{m}\right)+\mu y_{m}\left(\tau_{m}{ }^{\prime}, t_{0 m}, D_{m}, \mu\right)-x_{01}\left(\tau_{m}{ }^{\prime}-T, t_{01}, D_{1}\right)- \\
& \quad-\mu y_{1}\left(\tau_{m}{ }^{\prime}-T, t_{01}, D_{1}, \mu\right)=0  \tag{2.8}\\
& \varphi\left[x_{0 k}\left(\tau_{k}{ }^{\prime}, t_{0 k}, D_{k}\right)+\mu y_{k}\left(\tau_{k}{ }^{\prime}, t_{0 k}, D_{k}, \mu\right)\right]=0 \quad(k=1, \ldots, m) \tag{2.9}
\end{align*}
$$

It is not difficult to show that for sufficiently small values of $\mu$ the system (2.7) to (2.9) has solutions which are close to some of the solutions of (2.5), provided that the corresponding of the parameters $h_{i}$ in (2.5) satisfy $l$ conditions of the type (1.6). It may also be shown that if these values $h_{i}$ satisfy

$$
\frac{\partial\left(P_{1}, \ldots, P_{l}\right)}{\partial\left(h_{1}, \ldots, h_{l}\right)} \neq 0
$$

then the solutions of the system (2.7) to (2.9) are unique and correspond to a single solution of the system (1.1). On the other hand, this proof need not be carried out in full, since it follows immediately from the results of [1,2].
3. Equations of linear approximation. Let a set of values of the parameters $h_{1}, \ldots . h_{l}$ be given. At the same time, let there be given one of the solutions of Equations (1.3), i.e. all $C_{k}$ and $\tau_{k}$ are determined. In each of the domains $G_{k}$ one may choose an initial instant in the interval $\tau_{k-1} \leqslant t_{0 k} \leqslant \tau_{k}$, let us choose it such that

$$
\begin{equation*}
t_{0 k}=\tau_{k} \tag{3.1}
\end{equation*}
$$

Then, obviously, we shall have

$$
\begin{equation*}
x_{0 k}\left(\tau_{k}, \tau_{k}, C_{k}\right)=C_{k},\left\|\frac{\partial x_{0 k}\left(\boldsymbol{\tau}_{k}, \boldsymbol{\tau}_{k}, C_{k}\right)}{\partial C_{k}}\right\|=E_{n} \tag{3.2}
\end{equation*}
$$

where $E_{n}$ is the identity matrix of $n$ rows and $n$ columns.
It may be shown that the conditions (3.1) are not of use for substituting in (2.6), because it may happen that $t_{0 k}$ does not lie in the interval $\left[\tau_{k-1}{ }^{\prime}, \tau_{k}^{\prime}\right]$ and that the point $D_{k}$ is outside the domain $G_{k}$. On the other hand, it may readily be shown that the domain of definition of the functions $X$ and $f$ may be extended, preserving the conditions assuring the existence and uniqueness of solutions of the system (1.1),
so as to contain a certain neighborhood of the boundary of the domain $G_{k}$, and that the point $D_{k}$ will lie in this neighborhood when $\mu$ is sufficiently small.

We may now seek the solution of Equations (1.1) of the form (2.6), satisfying the conditions (2.7) to (2.9). If such a solution exists for arbitrarily chosen, sufficiently small $\mu$, then as $\mu \rightarrow 0$ it converges continuously to the chosen generating solution, and the system (2.7) to (2.9) must then possess a solution which is close to the solution of (2.2) to (2.4), that is, the $r_{k}^{\prime}$ and $D_{k}$ must differ but little from the $r_{k}$ and $C_{k}$. Consequently, with an accuracy to higher-order terms, the conditions (2.7) to (2.9) must be equivalent to the following relations:

$$
\begin{align*}
& x_{0 k}\left(\tau_{k}, \tau_{k}, C_{k}\right)+\left\|\frac{\partial x_{0 k}\left(\tau_{k}, \tau_{k}, C_{k}\right)}{\partial C_{k}}\right\| \delta C_{k}+\frac{\partial x_{0 k}}{} \frac{\left(\tau_{k}, \tau_{k}, C_{k}\right)}{\partial t} \delta \tau_{k}+  \tag{3.3}\\
& +\mu y_{k}\left(\tau_{k}, \tau_{k}, C_{k}, 0\right)-x_{0, k+1}\left(\tau_{k}, \tau_{k+1}, C_{k+1}\right)-\left\|\frac{\partial x_{0, k+1}\left(\tau_{k}, \tau_{k+1}, C_{k+1}\right)}{\partial C_{k+1}}\right\| \delta C_{k+1}- \\
& -\mu y_{k+1}\left(\tau_{k}, \tau_{k+1}, C_{k+1}, 0\right)-\frac{\partial x_{0, k+1}\left(\tau_{k}, \tau_{k+1}, C_{k+1}\right)}{\partial t} \delta \tau_{k}=0(k=1, \ldots, m-1) \\
& x_{0 m}\left(\tau_{m}, \tau_{m}, \quad C_{m}\right)+\left\|\frac{\partial x_{0 m}\left(\tau_{m}, \tau_{m}, C_{m}\right)}{\partial C_{m}}\right\| \delta C_{m}+\frac{\partial x_{0 m}\left(\tau_{m}, \tau_{m}, C_{m}\right)}{\partial t} \delta \tau_{m}+  \tag{3.4}\\
& \quad+\mu y_{m}\left(\tau_{m}, \tau_{m}, C_{m}, 0\right)-x_{01}\left(\tau_{m}-T, \tau_{1}, C_{1}\right)--\left\|\frac{\partial x_{01}\left(\tau_{m}-T_{1}, \tau_{1}, C_{1}\right)}{\partial C_{1}}\right\| \delta C_{1}- \\
& \quad-\frac{\partial x_{01}\left(\tau_{m}-T, \tau_{1}, C_{1}\right)}{\partial t} \delta \tau_{m}-\mu y_{1}\left(\tau_{m}-T, \tau_{1}, C_{1}, 0\right)=0 \\
& \varphi_{k}\left(x_{0 k}\right)+\frac{\partial \varphi_{k}}{\partial x_{0 k}}\left[\left\|\frac{\partial x_{0 k}\left(\tau_{k}, \tau_{k}, C_{k}\right)}{\partial C_{k}}\right\| \delta C_{k}+\frac{\partial x_{0 k}\left(\tau_{k}, \tau_{k}, C_{k}\right)}{\partial t} \delta \tau_{k}+\right.  \tag{3.5}\\
& \left.+\mu y_{k}\left(\tau_{k}, \tau_{k}, C_{k}, 0\right)\right]=0
\end{align*}
$$

where $\delta r_{k}=\tau_{k}^{\prime}-\tau_{k}, \delta C_{k}=D_{k}-C_{k}$. In view of (2.2) to (2.4) and (3.2), Equations (3.3). (3.4) and (3.5) may he simplified; using the notation $y_{0 k}(t)=y_{k}\left(t, r_{k}, C_{k}, 0\right)$, we obtain

$$
\begin{equation*}
\delta C_{k}-\left\|\frac{\partial x_{0, k+1}\left(\tau_{k}, \tau_{k+1}, C_{k+1}\right)}{\partial C_{k+1}}\right\| \delta C_{k+1}+\left[\frac{\partial x_{0 k}\left(\tau_{k}, \tau_{k}, C_{k}\right)}{\partial t}-\right. \tag{3.6}
\end{equation*}
$$

$\left.-\frac{\partial x_{0, k+1}\left(\tau_{k}, \tau_{k+1}, C_{k+1}\right)}{\partial t}\right] \delta \tau_{k}=\mu\left[y_{0, k+1}\left(\tau_{k}\right)-y_{0 k}\left(\tau_{k}\right)\right] \quad(k=1, \ldots, m-1)$
$\delta C_{m}-\left\|\frac{\partial x_{01}\left(\tau_{m}-T, \tau_{1}, C_{1}\right)}{\partial C_{1}}\right\| \delta C_{1}+\left[\frac{\partial x_{\mathrm{a} m}\left(\tau_{m}, \tau_{m}, C_{m}\right)}{\partial t}-\right.$

$$
\begin{equation*}
\left.-\frac{\partial x_{01}\left(\tau_{m}-T, \tau_{1}, C_{1}\right)}{\partial t}\right] \delta \tau_{m}=\mu\left[y_{01}\left(\tau_{m}-T\right)-y_{0 m}\left(\tau_{m}\right)\right] \tag{3.7}
\end{equation*}
$$

$$
\begin{equation*}
\frac{\partial \varphi_{k}}{\partial x_{0 k}}\left[\delta C_{k}+\frac{\partial x_{0 k}\left(\tau_{k}, \tau_{k}, C_{k}\right)}{\partial t} \delta \tau_{k}\right]=-\mu \frac{\partial \varphi_{k}}{\partial x_{0 k}} y_{0 k}\left(\tau_{k}\right) \quad(k=1, \ldots, m) \tag{3.8}
\end{equation*}
$$

The terms neglected here, as well as those neglected in previous formulas, are small of the second order, because the conditions imposed on the functions $X$ insure the continuity of the derivatives of the $x_{0 k}$ with respect to all its arguments; consequently, the values of these derivatives for $\delta \tau_{k}=\delta C_{k}=0$ and the values of these derivatives for certain mean values of the arguments (at which mean values of formulas written are exactly true) differ but little for sufficiently small $\mu$. Further, substituting from (2.6) into (1.1) and retaining only the firstorder terms in these equations, we arrive at

$$
\begin{equation*}
\frac{d y_{0 k}}{d t}=\left\|\frac{\partial X\left(x_{0 k}, t\right)}{\partial x_{0 k}}\right\| y_{0 k}(t)+f_{0 k}(t) \quad(k=1, \ldots, m) \tag{3.9}
\end{equation*}
$$

where

$$
f_{0 k}(t)=f\left[x_{0 k}\left(t, \tau_{k}, G_{k}\right), t, 0\right]
$$

In the terminology of [4], the differential equations (3.9), together with the conditions (3.6) to (3.8), are called the linear approximation to (1.1).

Equations (3.6) to (3.8) constitute a linear nonhomogeneous system for the determination of the unknowns $\delta C_{k}$ and $\delta \tau_{k}$. The coefficients of the corresponding homogeneous system form a matrix whose determinant coincides with the Jacobian of the system (2.2) to (2.4). Therefore this homogeneous system possesses $l$ independent solutions. In such a case, as is known, in order that there exist solutions of the nonhomogeneous system it is necessary and sufficient that the vector of order $m(n+1)$ formed by the right-hand terms of the nonhomogeneous system be orthogonal to all $l$ independent solutions of the adjoint homogeneous system (i.e. the homogeneous system whose matrix of coefficients is the transpose of the matrix of the coefficients of the original homogeneous system).
4. Equations of variation. Setting

$$
\begin{equation*}
x_{k}=x_{0 k}\left(t, \tau_{k}, C_{k}\right)+z_{k}(t) \quad(k=1, \ldots, m) \tag{4.1}
\end{equation*}
$$

and supposing that $z_{k}$ is small, we obtain the equations of variation for the periodic solutions of (1.3)

$$
\begin{equation*}
\frac{d z_{k}}{d t}=\left\|\frac{\partial X\left(x_{0 k}, t\right)}{\partial x_{0 k}}\right\| z_{k} \quad(k=1, \ldots, m) \tag{4.2}
\end{equation*}
$$

The matrix $\left\|\partial X / \partial x_{0 k}\right\|$ is continuous* in the domain $G_{k}$ and Equation (1.3) has the general solution (2.1), hence Equation (4.2) has, in this domain, the general solution

$$
\begin{equation*}
z_{k}=\left\|\frac{\partial x_{0 k}\left(t, \tau_{k}, C_{k}\right)}{\partial C_{k}}\right\| A_{k} \quad(k=1, \ldots, m) \tag{4.3}
\end{equation*}
$$

where $A_{k}$ is an arbitrary column vector, which is small, in view of the smallness of the functions $z_{k}$.

Let us now explain how the constants $A_{k}$ may be chosen in order that the variational motion (4.1) satisfy the continuity condition on the boundary of the domains $G_{k}$ and the condition of periodicity with period T. Substituting from (4.3) into (4.1), and then setting (4.1) into (2.2) to (2.4), we obtain

$$
\begin{align*}
& x_{0 k}\left(\tau_{k}{ }^{*}, \tau_{k}, C_{k}\right)+\left\|\frac{\partial x_{0 k}\left(\tau_{k}^{*}, \tau_{k}, C_{k}\right)}{\partial C_{k}}\right\| A_{k}-x_{0, k+1}\left(\tau_{k}^{*}, \tau_{k+1}, C_{k+1}\right)-  \tag{4.4}\\
& -\left\|\frac{\partial x_{0, k+1}\left(\tau_{k}^{*}, \tau_{k+1}, C_{k+1}\right)}{\partial C_{k+1}}\right\| A_{k+1}=0 \quad(k=1, \ldots, m-1) \\
& x_{0 m}\left(\tau_{m}^{*}, \tau_{m}, C_{m}\right)+\left\|\frac{\partial x_{0 m}\left(\tau_{m}^{*}, \tau_{m}, C_{m}\right)}{\partial C_{m}}\right\| A_{m}-x_{01}\left(\tau_{m}^{*}-T, \tau_{1}, C_{1}\right)-  \tag{4.5}\\
& -\left\|\frac{\partial x_{01}\left(\tau_{m}^{*}-T, \tau_{1}, C_{1}\right)}{\partial C_{1}}\right\| A_{1}=0 \\
& \varphi_{k}\left[x_{0 k}\left(\tau_{k}^{*}, \tau_{k}, C_{k}\right)+\left\|\frac{\partial x_{0 k}\left(\tau_{k}^{*}, \tau_{k}, C_{k}\right)}{\partial C_{k}}\right\| A_{k}\right]=0 \quad(k=1, \ldots, m) \tag{4.6}
\end{align*}
$$

Here $r_{k}{ }^{*}$ is that instant of time at which the variational integral curve (4.1) intersects the surface of discontinuity.

The system of equations (4.4) to (4.6) differs from the system (2.2) to (2.4) by terms of small order. Consequently, $\tau_{k}{ }^{*}$ differs slightly from $r_{k}$ and, up to higher-order terms, the system (4.4) to (4.6) is equivalent to the system

$$
\begin{align*}
& \left\|\frac{\partial x_{0 k}}{\partial C_{k}}\right\| A_{k}-\left\|\frac{\partial x_{0, k+1}^{\prime}}{\partial C_{k+1}}\right\| A_{k+1}+\left[\frac{\partial x_{0 k}\left(\tau_{k}\right)}{\partial t}-\frac{\partial x_{0, k+1}^{\prime}\left(\tau_{k}\right)}{\partial t}\right] \Delta \tau_{k}=0  \tag{4.7}\\
& \left\|\frac{\partial x_{0 m}}{\partial C_{m}}\right\| A_{m}-\left\|\frac{\partial x_{01}^{\prime}}{\partial C_{1}}\right\| A_{1}+\left[\frac{\partial x_{0 m}\left(\tau_{m}\right)}{\partial t}-\frac{\left.\partial x_{01^{\prime}\left(\tau_{m}-T\right)}^{\partial t}\right] \Delta \tau_{m}=0}{(k=1, \ldots}\right. \tag{4.8}
\end{align*}
$$

* In all matrices such as this one, each row consists of the partial derivatives of one and the same function with respect to the independent variables in question.

$$
\begin{equation*}
\frac{\partial \varphi_{k}}{\partial x_{0 k}}\left[\left\|\frac{\partial x_{0 k}}{\partial C_{k}}\right\|_{A_{k}}+\frac{\partial x_{0 k}\left(\tau_{k}\right)}{\partial t} \Delta \tau_{k}\right]=0 \quad(k=1, \ldots m) \tag{4.9}
\end{equation*}
$$

where

$$
\begin{aligned}
& \left\|\begin{array}{l}
\frac{\partial x_{0 k}}{\partial C_{k}} \|
\end{array}\right\|=\left\|\frac{\partial x_{0 k}\left(\tau_{k}, \tau_{k}, C_{k}\right)}{\partial C_{k}}\right\|, \quad\left\|\frac{\partial x_{0, k+1}}{\partial C_{k+1}}\right\| \equiv\left\|\frac{\partial x_{0, k+1}\left(\tau_{k}, \tau_{k+1}, C_{k+1}\right)}{\partial C_{k+1}}\right\| \\
& \left\|\frac{\partial x_{01}}{\partial C_{1}}\right\| \equiv\left\|\frac{\partial x_{01}\left(\tau_{m}-T, \tau_{1}, C_{1}\right)}{\partial C_{1}}\right\|, \quad \frac{\partial x_{0 k}\left(\tau_{k}\right)}{\partial t} \equiv \frac{\partial x_{0 k}\left(\tau_{k}, \tau_{k}, C_{k}\right)}{\partial t} \\
& \\
& \\
& \frac{\partial x_{0, k+1}\left(\tau_{k}\right)}{\partial t}=\frac{\partial x_{0, k+1}\left(\tau_{k}, \tau_{k+1}, C_{k+1}\right)}{\partial t}, \quad \Delta \tau_{k}=\tau_{k}{ }^{*}-\tau_{k}
\end{aligned}
$$

Equations (4.9) define uniquely the $\Delta r_{k}$ :

$$
\begin{equation*}
\Delta \tau_{k}=-\frac{1}{d \varphi_{k}\left(\tau_{k}\right) d t}\left(\left\|\frac{\partial x_{0 k}}{\partial C_{k}}\right\| A_{k}\right) \frac{\partial \varphi}{\partial x_{0 k}} \quad\left(\frac{d \varphi_{k}\left(\tau_{k}\right)}{d t} \neq 0 \text { in view of (1.5) }\right) \tag{4.10}
\end{equation*}
$$

From Equation (1.3) we obtain

$$
\begin{equation*}
\frac{\partial x_{0 . k}\left(\tau_{k}\right)}{\partial t}-\frac{\partial x_{0, k+1}\left(\tau_{k}\right)}{\partial t}=X_{k}\left(\tau_{k}\right)-X_{k+1}\left(\tau_{k}\right)=-\Delta_{k} \tag{4.11}
\end{equation*}
$$

where $\Delta_{k}$ denotes the fump of the function $X$ across the surface of discontinuity.

Making use of (4.10) and (4.11) in (4.7) and (4.8), keeping in mind (3.2), we obtain

$$
\begin{equation*}
A_{k}-\left\|\frac{\partial x_{0, k+1}^{\prime}}{\partial C_{k+1}}\right\| A_{k+1}+\Delta_{k} \frac{A_{k} \partial \varphi_{k} / \partial x_{0 k}}{d \varphi_{k}\left(\tau_{k}\right) / d t}=0 \quad\binom{k=1, \ldots, m}{m+1 \equiv 1} \tag{4.12}
\end{equation*}
$$

The system of equations* which has been obtained is equivalent to the homogeneous system which corresponds to (3.6) to (3.8). Therefore it also must have $l$ independent solutions, which determine $l$ continuous periodic solutions of the variational equations (4.1). It should be observed that the functions $z_{k}$ themselves need not, in general, be continuous. It is only when all $\Delta_{k}$ are equal to zero, that is, when $X$ is continuous on every periodic solution, that the conditions (4.12) coincide with the requirement that the functions $z_{k}$ be continuous.

* It is easy to see that (4.12), together with (4.2), coincide with the linear approximation of the equations (1.3) which is obtained in [4].

Thus we may construct in this manner a set of $l$ periodic solutions of the equations of variation (4.2), satisfying the conditions (4.12).

Let us now consider the system of linear equations which is adjoint to (4.2)

$$
\begin{equation*}
\frac{d u_{k}}{d t}+\left\|\frac{\partial X\left(x_{0 k}, t\right)}{\partial x_{0 k}}\right\|^{*} u_{k}=0 \quad(k=1, \ldots, m) \tag{4.13}
\end{equation*}
$$

and the system of linear equations

$$
\begin{equation*}
B_{k+1}-\left\|\frac{\partial x_{0, k+1}^{\prime}}{\partial C_{k+1}}\right\|^{*} B_{k}+\frac{\partial \varphi_{k+1}\left(\tau_{k+1}\right)}{\partial x_{0 k}} \frac{\left(\Delta_{k+1}\left(\tau_{k+1}\right) B_{k+1}\right)}{d \varphi_{k+1}\left(\tau_{k+1}\right) / d t}=0 \tag{4.14}
\end{equation*}
$$

Here, and in the following, the asterisk denotes the transposed matrix.

It is readily verified that the matrix of coefficients of the system (4.14) is the transpose of the matrix of coefficients of (4.12). Hence the system (4.14) also possesses $l$ independent solutions

$$
\begin{equation*}
B_{k}^{(1)}, \quad B_{k}^{(2)}, \ldots, \quad B_{k}^{(l)} \quad(k=1, \ldots, m) \tag{4.15}
\end{equation*}
$$

Let us now seek a solution of (4.13) satisfying the conditions

$$
\begin{equation*}
u_{k+1}^{(i)}\left(\tau_{k}\right)=B_{k}^{(i)} \tag{4.16}
\end{equation*}
$$

Inserting (4.16) into (4.14), we obtain

$$
\begin{gather*}
u_{k+2}^{(i)}\left(\tau_{k+1}\right)-\left\|\frac{\partial x_{0, k+1}^{\prime}}{\partial C_{k+1}}\right\|_{k+1}^{*}\left(\tau_{k}\right)+\frac{\partial \varphi_{k+1}\left(\tau_{k+1}\right)}{\partial x_{0 k}} \frac{\left(\Delta_{k+1}\left(\tau_{k+1}\right) u_{k+2}^{(i)}\left(\tau_{k+1}\right)\right)}{d \varphi_{k+1}\left(\tau_{k+1}\right) / d t}=0 \\
\left(k=0,1, \ldots, m-1 ; \tau_{0}=\tau_{m}-T\right) \tag{4.17}
\end{gather*}
$$

In view of well-known properties of the solutions of adjoint systems, we obtain

$$
\begin{equation*}
\left\|\frac{\partial x_{0, k+1}^{\prime}}{\partial C_{k+1}}\right\|_{k+1}{ }^{(i)}\left(\tau_{k}\right)=\left\|\frac{\partial x_{0, k+1}}{\partial C_{k+1}}\right\|_{i}^{*} u_{k+1}{ }^{(i)}\left(\tau_{k+1}\right)=u_{k+1}{ }^{(i)}\left(\tau_{k+1}\right) \tag{4.18}
\end{equation*}
$$

because

$$
\left\|\frac{\partial x_{0, k+1}}{\partial C_{k+1}}\right\|^{*}=E_{n}^{\prime}
$$

Equations (4.17) then become

$$
\begin{equation*}
u_{k+1}{ }^{(i)}\left(\tau_{k}\right)-u_{k}^{(i)}\left(\tau_{k}\right)+\frac{\partial \varphi_{k}}{\partial x_{0 k}}\left(\tau_{k}\right) \frac{\left(\Delta_{k}\left(\tau_{k}\right) \cdot u_{k+1}{ }^{(i)}\left(\tau_{k}\right)\right)}{d \varphi_{k}\left(\tau_{k}\right) / d t}=0 \quad(k=1, \ldots, m) \tag{4,19}
\end{equation*}
$$

Obviously, if the function $X$ is continuous (i.e. all the $\Delta_{k}=0$ ) then (4.19) defines a periodic solution of the system of equations (4.13).
5. Conditions for the existence of a periodic solution. Let us return to the solution of Equation (3.9). It is readily seen that this equation has as a solution (for $t \leqslant \tau_{k}$ )

$$
\begin{equation*}
y_{0 k}=\int_{\tau_{k}}^{t}\left\|\frac{\partial x_{0 k}\left(t, \tau, C_{k}\right)}{\partial C_{k}}\right\| f_{0 k}(\tau) d \tau \quad(k=1, \ldots, m) \tag{5.1}
\end{equation*}
$$

provided that it is required that

$$
\begin{equation*}
\left\|\frac{\partial x_{0 k}\left(\tau, \tau, C_{k}\right)}{\partial C_{k}}\right\|=E_{n} \tag{5.2}
\end{equation*}
$$

Inserting (5.1) into the conditions (3.6) to (3.8), we obtain the following system:

$$
\begin{gather*}
\delta C_{k}-\left\|\frac{\partial x_{0, k+1}^{\prime}}{\partial C_{k+1}}\right\| \delta C_{k+1}-\Delta_{k} \delta \tau_{k}=\mu \int_{\tau_{k+1}}^{\tau_{k}} \| \frac{\partial x_{0, k+1}\left(\tau_{k}, \tau, C_{k+1}\right)}{\partial C_{k+1}} f_{0, k+1}(\tau) d \tau  \tag{5.3}\\
\bar{\partial} C_{m}\left\|\frac{\partial x_{01}}{\partial C_{1}}\right\| \delta C_{1}-\Lambda_{m} \delta \tau_{m}=\mu \int_{\tau_{1}}^{\tau}\left\|\frac{\partial x_{01}\left(\tau_{m}-T, \tau, C_{1}\right)}{\partial C_{1}}\right\| f_{01}(\tau) d \tau \\
\frac{\partial \varphi_{k}}{\partial x_{0 k}}\left[\delta C_{k}+\frac{\partial x_{0 k}}{\partial t} \tau_{k}\right)  \tag{5.4}\\
\left.\partial \tau_{k}\right]=0 \quad(k=1, \ldots, m-1) \tag{5.5}
\end{gather*}
$$

Taking into account that (4.16) is a solution of the corresponding homogeneous system, the conditions for the existence of the system (5.3) to (5.5) may be written thus:

$$
\begin{align*}
& \quad \sum_{k=1}^{m-1} \int_{\tau_{k+1}}^{\tau_{k}}\left\|\frac{\partial x_{0, k+1}\left(\tau_{k}, \tau, C_{k+1}\right)}{\partial C_{k+1}}\right\| f_{0, k+1}(\tau) u_{k+1}{ }^{(i)}\left(\tau_{k}\right) d \tau+ \\
& +\int_{-1}^{=m^{-T}}\left\|\frac{\partial x_{01}\left(\tau_{m}-T, \tau, C_{1}\right)}{\partial C_{1}}\right\| f_{01}(\tau) u_{1}^{(i)}\left(\tau_{m}-T\right) d \tau=0 \quad(i=1, \ldots, l) \tag{5.6}
\end{align*}
$$

Making use of (4.18), and denoting by $f_{0}(t)$ the function which equals $f_{0 k}$ on each domain $G_{k}$, and denoting by $u^{(i)}(t)$ the function which equals $u_{k}{ }^{(i)}(t)$ in $G_{k}$ we obtain the conditions for the existence of a periodic solution in the final form:

$$
\begin{equation*}
\int_{\tau_{m}-T}^{\tau_{m}} f_{0}(\tau) u^{(i)}(\tau) d \tau=0 \quad(i=1, \ldots, l) \tag{5.7}
\end{equation*}
$$

For the determination of the parameters $h_{1}, \ldots, h_{l}$ we obtain the system

$$
\begin{equation*}
P_{i}\left(h_{1}, \ldots, h_{l}\right) \equiv \int_{0}^{T} f_{0}(\tau) u^{(i)}(\tau) d \tau=0 \quad(i=1, \ldots, l) \tag{5.8}
\end{equation*}
$$

If the system (5.8) has a unique solution $h_{1}{ }^{*}, \ldots, h_{l}{ }^{*}$, i.e. if the Jacobian

$$
\left[\frac{\partial\left(P_{1}, \ldots, P_{l}\right)}{\partial\left(h_{1}, \ldots, h_{l}\right)}\right]_{h_{j}=h_{j}^{*}} \neq 0
$$

then the system (1.1) will have a unique solution for all sufficiently small values of the parameter $\mu$, a solution which is close to the solution of the generating system.

The conditions (5.8) coincide formally with the conditions obtained in [3] for the case of equations whose right-hand sides possess continuous partial derivatives of the second order.

However, the functions $u^{(i)}$ need not be continuous in the present case: they must satisfy the conditions (4.19), which in the present case are

$$
\begin{gather*}
u_{s, k+1}^{(i)}\left(\tau_{k}\right)-u_{s, k}^{(i)}\left(\tau_{k}\right)+\frac{\partial \varphi_{k}}{\partial x_{0 s}}\left(\tau_{k}\right) \frac{1}{d \varphi_{k}\left(\tau_{k}\right) / d t} \sum_{j=1}^{n} \Delta_{k j}\left(\tau_{k}\right) u_{j, k+1}^{(i)}\left(\tau_{k}\right)=0  \tag{5.9}\\
(i=1, \ldots, l, s=1, \ldots, n ; k=1, \ldots, m)
\end{gather*}
$$

Note. Conditions (5.7) and (5.9) still hold when the surface of discontinuity is given in terms of a perindic function (with period $T$ ) of the time

$$
\begin{equation*}
\varphi_{k}(x, t)=0 \tag{5.10}
\end{equation*}
$$

In this case, condition (1.5) must be replaced by

$$
\begin{equation*}
\frac{d \varphi_{k}}{d t}=\frac{\partial \varphi_{k}}{\partial x} \frac{\partial x_{0}}{\partial t}+\frac{\partial \varphi_{k}}{\partial t}=0 \tag{5.11}
\end{equation*}
$$

6. Quasiconservative systems. By way of an example, illustrating the preceding results, consider a system which is "close" to a conservative system

$$
\begin{align*}
& \dot{q}_{s}=\frac{\partial H}{\partial p_{s}}+\mu Q_{8}\left(q_{1}, \ldots, q_{n}, p_{1}, \ldots, p_{n}, t, \mu\right) \\
& \dot{p}_{\mathrm{s}}=-\frac{\partial H}{\partial q_{s}}+\mu P_{s}\left(q_{1}, \ldots, q_{n}, p_{1}, \ldots, p_{n}, t, \mu\right)
\end{align*}
$$

where $H$ is the Hamiltonian of the generating conservative system, supposed not to depend explicitly upon $t$, and $Q_{s}$ and $P_{s}$ are periodic functions of $t$ with period $T$.

Let us assume that the forces acting in the generating system

$$
\begin{equation*}
\dot{q}_{s 0}=\frac{\partial H}{\partial p_{s 0}}, \quad \dot{p}_{s 0}=-\frac{\partial H}{\partial q_{s 0}} \tag{6.2}
\end{equation*}
$$

are conservative forces which are discontinuous with respect to the coordinate $q_{k}$, with a discontinuity of the first kind on the surfaces

$$
\begin{equation*}
\varphi_{k}\left(q_{1}, \ldots, q_{n}\right)=0 \tag{6.3}
\end{equation*}
$$

In this case the derivatives $\partial H / \partial q_{s}$ also have discontinuities of the first kind on these same surfaces.

Suppose that the system (6.2) has a periodic solution with period $T$. Then, in view of the explicit independence of $H$ upon $t$, it also has the one parameter family of solutions

$$
\begin{equation*}
q_{s 0}=q_{s 0}(t+h), \quad p_{s 0}=p_{s 0}(t+h) \tag{6.4}
\end{equation*}
$$

The adjoint system to the equations of variation

$$
\begin{equation*}
\dot{u}_{s}=-\sum_{j} \frac{\partial^{2} H}{\partial p_{j} \partial q_{s}} u_{j}+\sum_{j} \frac{\partial^{2} H}{\partial q_{j} \partial q_{s}} v_{j}, \quad \dot{v}_{s}=-\sum_{j} \frac{\partial^{2} H}{\partial p_{j} \partial p_{s}} u_{j}+\sum_{j} \frac{\partial^{2} H}{\partial q_{j} \partial p_{s}} v_{j} \tag{6.5}
\end{equation*}
$$

possesses a family of periodic solutions

$$
\begin{equation*}
u_{s}=-\dot{p}_{s 0}(t+h), \quad v_{s}=\dot{q}_{s 0}(t+h) . \tag{6.6}
\end{equation*}
$$

as may be readily verified by direct substitution.
Let us show that the solutions (6.6) do satisfy conditions (5.9). Indeed, the equations for $v_{s}$ are identically fulfilled, since

$$
v_{s, k+1}-v_{s, k}=0, \quad \frac{\partial \varphi_{k}}{\partial p_{s}}=0
$$

The equations for the $u_{s}$ are

$$
\sum_{j=1}^{2 n} \Delta_{k j} u_{j, k+1}=\sum_{j=1}^{n}\left[\left(\frac{\partial H}{\partial q_{j 0}}\right)_{k}-\left(\frac{\partial H}{\partial q_{j 0}}\right)_{k+1}\right] \dot{q}_{j 0}=\left(\frac{d \Pi}{d t}\right)_{k}-\left(\frac{d \Pi}{d t}\right)_{k+1}
$$

where $\Pi$ is the potential energy of the system; hence

$$
\begin{aligned}
& u_{s, k+1}-u_{s, k}+\frac{\partial \varphi_{k}}{\partial q_{s 0}} \frac{1}{d \varphi_{k} / d t} \sum_{j=1}^{n} \Delta_{k j} u_{j, k+1} \\
& =\dot{p}_{s 0, k}-\dot{p}_{s 0, k+1}+\frac{\partial \varphi_{k}}{\partial q_{s 0}} \frac{1}{d \varphi_{k} / d t}\left[\left(\frac{d \Pi}{d t}\right)_{k}-\left(\frac{d \Pi}{d t}\right)_{k+1}\right] \\
& \quad=\left(\frac{\partial H}{\partial q_{s 0}}\right)_{k+1}-\left(\frac{\partial H}{\partial q_{s 0}}\right)_{k}+\left(\frac{\partial \Pi}{\partial q_{s 0}}\right)_{k}-\left(\frac{\partial \Pi}{\partial q_{s 0}}\right)_{k+1}=0
\end{aligned}
$$

because in the case under consideration the jump of $\partial H / \partial q_{s 0}$ equals the jump of $\partial \Pi / \partial q_{s} 0^{\circ}$

Finally, therefore, the condition for the existence of a periodic solution of the system which is close to a solution of (6.4), may be written
$\int_{0}^{T} \sum_{s}\left[P_{s}\left(q_{j 0}, p_{j 0}, \tau, 0\right) \dot{q}_{s 0}(\tau+h)-Q_{s}\left(q_{j 0}, p_{j 0}, \tau, 0\right) \dot{p}_{s 0}(\tau+h)\right] d \tau=0$
For a second-order system

$$
\begin{equation*}
\ddot{x}+F(x)=\mu f(x, \dot{x}, t, \mu) \tag{6.8}
\end{equation*}
$$

the condition (6.7) assumes the known form

$$
\begin{equation*}
\int_{0}^{T} f\left(x_{0}, \dot{x}_{0}, \tau, 0\right) \dot{x}_{0}(\tau+h) d \tau=0 \tag{6.9}
\end{equation*}
$$

where $x_{0}(t+h)$ is the family of periodic solutions of the generating system

$$
\ddot{x}_{0}+F\left(x_{0}\right)=0
$$

Condition (6.9) was obtained earlier for analytic $F(x)$ and $f(x, x$, $t, \mu$ ) (see e.g. [5]); it remains valid when these functions possess a finite number of discontinuities of the first kind in the domain in question.

## BIBLIOGRAPHY

1. Neimark, Iu.I., Metod tochechnykh otobrazhenii v teorii nelineinykh kolebanii (Method of point transformations in the theory of nonlinear oscillations). I, II, III. Izv. Min. vyssh. obraz. Radiofiz. Nos. 1 and 2, 1958.
2. Neimark, Iu. I. and Shil'nikov, L.P., o primenenii metoda malogo parametra k sistemam differentsial' nykh uravnenii s razryvnymi pravymi chastiami (on the application of the method of small parameter to systems of differential equations with discontinuous right-hand sides). Izv. Akad. Nauk SSSR, OTN, Mekhanika i mashinostroenie, No. 6, 1959.
3. Malkin, I. G. Nekotorye zadachi teorii nelineinykh kolebanii (Some Problems of the Theory of Nonlinear Oscillations). Gostekhizdat, 1956.
4. Aizerman, M. A. and Gantmakher, F. R., Ustoichivost' po lineinomy priblizheniiu periodicheskovo resheniia sistemy uravnenii s razryvnymi pravymi chastiami (Stability in linear approximation of periodic solution of a system of equations with discontinuous right-hand sides). $P M M$ Vol. 21, No. 5, 1957.
5. Kats, A. M., Vynushdennye kolebaniia nelineinykh sistem s odnot stepen'iu svobody, blizkikh k conservativnym (Forced oscillations of nonlinear systems with one degree of freedom which are quasiconservative). $P M M$ Vol. 19, No. $1,1955$.
